

NUMERICAL ANALYSIS OF FRACTIONAL ORDER INTEGRAL DYNAMICAL MODELS WITH PIECEWISE CONTINUOUS KERNELS*A. Tynda*¹, *D. Sidorov*^{2,3}, *I. Muftahov*^{2,4}¹Penza State University, Penza, Russian Federation²Melentiev Energy Systems Institute SB RAS, Irkutsk, Russian Federation³Irkutsk National Research Technical University, Irkutsk, Russian Federation⁴Main Computing Center of JSC Russian Railways, Irkutsk, Russian Federation

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Volterra integral equations find their application in many areas, including mathematical physics, control theory, mechanics, electrical engineering, and in various industries. In particular, dynamic Volterra models with discontinuous kernels are effectively used in power engineering to determine the operating modes of energy storage devices, as well as to solve the problem of load balancing. This article proposes the numerical scheme for solution of the fractional order linear Volterra integral equations of the first kind with piecewise continuous kernels. The developed approach is based on a polynomial collocation method and effectively approximate such a weakly singular integrals. The efficiency of proposed numerical scheme is illustrated by two examples.

Keywords: Volterra integral equations; numerical method; convergence; discontinuous kernel; singularity; fractional integral.

Introduction

Volterra integral equations (VIEs) [1] were introduced by Vito Volterra and then studied by Traian Lalescu in his 1908 thesis. In weakly regular case, VIEs were introduced in [2], and the theory of such equations is generalized to the case of systems of equations and to abstract operator equations in the monograph [3]. Integral equations are essential tools in various fields including power and electrical engineering, mathematical physics, control theory and mechanics.

Fractional integration and fractional differentiation are generalisations of notions of integer-order integration and differentiation, and include n -th derivatives and n -folded integrals (n denotes an integer number) as particular cases. The first mention of derivatives of non-integer order is presented in the correspondence between G. Leibniz and J. Bernoulli. In a letter written by G. Leibniz to G. L'opital, there is also an interesting mention about the paradox and possible useful practical application of differentials of order $\frac{1}{2}$. In recent years, interest in fractional calculus grows [4–6].

Fractional derivatives have numerous applications; they are in the core of various mathematical models for viscoelastic bodies, in chemical physics, the theory of gravity, viscoplasticity, etc. [6, 7]. Fractional derivatives have practical applications in modelling the behavior of viscoplastic materials. In particular, the equations of state in the theory of viscoplasticity contain fractional derivatives and can be reduced to weakly singular Volterra integral equations of the second kind [8, 9].

This article discusses weakly regular Volterra equations of the first kind

$$\int_0^t K(t, s)x(s)ds = f(t), \quad 0 \leq s \leq t \leq T, \quad f(0) = 0, \quad (1)$$

where the kernel is defined as follows:

$$K(t, s) = \begin{cases} K_1(t, s), & t, s \in m_1, & m_i = \{t, s \mid \alpha_{i-1}(t) < s < \alpha_i(t)\}, \\ \dots & \dots\dots\dots \\ K_n(t, s), & t, s \in m_n, & \alpha_0(t) = 0, \alpha_n(t) = t, i = \overline{1, n}, \end{cases} \quad (2)$$

$\alpha_i(t), f(t) \in C^1_{[0, T]}$, $K_i(t, s)$ have continuous derivatives with respect to t for $t \in \overline{m_i}$, $K_n(t, t) \neq 0, \alpha_i(0) = 0, 0 < \alpha_1(t) < \alpha_2(t) < \dots < \alpha_{n-1}(t) < t, \alpha_1(t), \dots, \alpha_{n-1}(t)$ grow in a small neighborhood $0 \leq t \leq \tau, 0 < \alpha'_1(0) \leq \dots \leq \alpha'_{n-1}(0) < 1$.

The works [2, 10] are the first to study such equations and to outline the non-uniqueness of the solution. The paper [11] derives the existence of a continuous solution depending on free parameters as well as sufficient conditions for the existence of a unique continuous solution to the system of VIE of the first kind with discontinuous kernels. The class of Volterra operator equations of the first kind with piecewise continuous kernels is studied in [12]. Various numerical methods for linear and nonlinear VIE with piecewise continuous kernels and their systems with applications for power systems operation are proposed in [13–15]. In [16], the Volterra model is employed for a load leveling problem in modelling the hybrid AC/DC power systems with renewable energy sources and storage system.

For systematic studies of VIE of the first kind with piecewise continuous kernels, see the book [3] and the part 1 in [17]. For numerical solution of Volterra integral equation of Abel type, see [18]. The numerical solution of the Volterra fractional integral equations of the second kind using the Simpson 3/8 rule method is proposed in [4].

In the case of a fractional order of integration for these equations, consider the left-sided Riemann–Liouville fractional integral of the order $\beta \geq 0$

$$I_t^\beta f(t) = \frac{1}{\Gamma(\beta)} \int_0^t \frac{f(s)}{(t-s)^\beta} ds. \quad (3)$$

For equation (1) in the case of fractional order of integration (3), we solve the integral equation

$$\frac{1}{\Gamma(\beta)} \int_0^t \frac{K(t, s)}{(t-s)^\beta} x(s) ds = f(t), \quad 0 \leq s \leq t \leq T, \quad f(0) = 0, \quad (4)$$

with piecewise continuous kernel (2).

1. Polynomial Collocation Method

Let us write equation (4) in the following expanded form:

$$\frac{1}{\Gamma(\beta)} \sum_{i=1}^n \int_{\alpha_{i-1}(t)}^{\alpha_i(t)} \frac{K_i(t, s)x(s)}{(t-s)^\beta} ds = f(t), \quad t \in [0, T], \quad 0 < \beta < 1. \quad (5)$$

Apply the following representation (as power series segment) of the solution to equation (5):

$$X_m(t) = \sum_{j=0}^m a_j t^j \tag{6}$$

with the desired $a_j, j = 0, 1, \dots, m$. Let us substitute (6) into (5), then we get

$$\frac{1}{\Gamma(\beta)} \sum_{i=1}^n \int_{\alpha_{i-1}(t)}^{\alpha_i(t)} \frac{K_i(t, s)}{(t-s)^\beta} \left(\sum_{j=0}^m a_j s^j \right) ds = f(t).$$

Let us change the order of integration and summation

$$\frac{1}{\Gamma(\beta)} \sum_{i=1}^n \sum_{j=0}^m a_j \int_{\alpha_{i-1}(t)}^{\alpha_i(t)} \frac{K_i(t, s) s^j}{(t-s)^\beta} ds = f(t) \tag{7}$$

and transform the last equality as follows:

$$\frac{1}{\Gamma(\beta)} \sum_{j=0}^m a_j \sum_{i=1}^n \int_{\alpha_{i-1}(t)}^{\alpha_i(t)} \frac{K_i(t, s) s^j}{(t-s)^\beta} ds = f(t).$$

We require the fulfillment of the last equality at the points of a uniform grid

$$t_k = \frac{kT}{m+1}, \quad k = 1, 2, \dots, m+1.$$

Therefore, we arrive at the following system of equations for the unknown coefficients $a_j, j = 0, 1, \dots, m$:

$$\sum_{j=0}^m \gamma_{kj} a_j = F_k, \quad k = 1, 2, \dots, m+1, \tag{8}$$

where

$$\gamma_{kj} = \frac{1}{\Gamma(\beta)} \sum_{i=1}^n \int_{\alpha_{i-1}(t_k)}^{\alpha_i(t_k)} \frac{K_i(t_k, s) s^j}{(t_k-s)^\beta} ds, \quad F_k = f(t_k). \tag{9}$$

System of linear algebraic equations (8) is solved for unknown expansion coefficients (6) by the Jordan–Gauss method, and in order to calculate the elements of the matrix of this system, below we construct special quadrature formulas that take into account the power singularity of the integrands.

1.1. Approximation of Weakly Singular Integrals

In the previous section, when constructing the matrix of system of linear algebraic equations, in order to determine the coefficients of the approximating polynomial, it is necessary to calculate the values of the functions

$$I(i, k, j) = \int_{\alpha_{i-1}(t_k)}^{\alpha_i(t_k)} \frac{K_i(t_k, s) s^j}{(t_k-s)^\beta} ds, \quad i = \overline{1, n}, \quad k = \overline{1, m+1}, \quad j = \overline{0, m}, \tag{10}$$

containing either a convergent improper integral (for $i = n$), or significant in magnitude integrands (for $i < n$). In this section, we propose a method to calculate integrals (10).

For each triple i, j, k , an internal mesh of nodes is constructed as follows:

$$\tau_l = \alpha_{i-1}(t_k) + \frac{l}{r} (\alpha_i(t_k) - \alpha_{i-1}(t_k)), \quad l = 0, 1, \dots, r, \quad r = \left\lceil \frac{\alpha_i(t_k) - \alpha_{i-1}(t_k)}{h} \right\rceil + 1, \quad (11)$$

where $[t]$ stands for integer part of a real number t , and h is a sufficiently small parameter, which we conventionally call the grid step. Then we approximate integral (10) as follows:

$$\begin{aligned} I(i, k, j) &= \sum_{l=1}^r \int_{\tau_{l-1}}^{\tau_l} \frac{K_i(t_k, s) s^j}{(t_k - s)^\beta} ds \approx \\ &\approx \frac{1}{1 - \beta} \sum_{l=1}^r K_i \left(t_k, \frac{\tau_{l-1} + \tau_l}{2} \right) \left(\frac{\tau_{l-1} + \tau_l}{2} \right)^j \left((t_k - \tau_{l-1})^{1-\beta} - (t_k - \tau_l)^{1-\beta} \right). \end{aligned} \quad (12)$$

Therefore, for relatively small values T , we propose a sufficiently effective method for solving equations of the form (5), the error of which can be estimated by the inequality

$$\|x(t) - X_m(t)\|_{C[0, T]} \leq \frac{L_{m+1} T^{m+1}}{(m+1)!}, \quad L = \max_{t \in [0, T]} |x^{(m+1)}(t)| \quad (13)$$

for $x(t) \in C^{m+1}[0, T]$.

With an increase in the length of the planning interval $[0, T]$, to maintain the order of accuracy, it is necessary to increase the degrees of the approximating polynomials, which leads to significant computational errors. In order to overcome this restriction, the method can be generalized using a polynomial spline approximation of the solution constructed in each section in a similar way.

It is known that exact solutions to weakly singular Volterra integral equations can have unbounded derivatives as $t \rightarrow 0$ [19]. In this case, the exact solutions belong to the class of functions $C^{r, \beta}(0, T]$ defined as follows.

Definition 11. Let $\Omega = (0, T]$, $0 \leq \beta < 1$. We say that a function $x(t)$ belongs to $C^{r, \beta}(0, T]$, if for $t \in (0, T]$ the function $x(t)$ has continuous derivatives up to the order r estimated as follows:

$$|x^{(k)}(t)| \leq \frac{A_k}{t^{k-1+\beta}} \quad \text{for all } t \in (0, T], k = 0, 1, \dots, r.$$

In this case, an approximate solution to equation (5) can be found as

$$X_m(t) = a_0 + \sum_{j=1}^m a_j t^{j-\beta}. \quad (14)$$

The method to determine coefficients expansion (14) is employed below in the similar way.

Table 1

Numerical results for Example 1

Errors				
m	$h = 10^{-2}$	$h = 10^{-3}$	$h = 10^{-4}$	$h = 10^{-5}$
3	$8,3730 \cdot 10^{-4}$	$1,0826 \cdot 10^{-3}$	$1,1168 \cdot 10^{-3}$	$1,0538 \cdot 10^{-3}$
4	$1,1624 \cdot 10^{-3}$	$2,0991 \cdot 10^{-4}$	$1,7066 \cdot 10^{-4}$	$1,6319 \cdot 10^{-4}$
5	$2,0042 \cdot 10^{-3}$	$5,3862 \cdot 10^{-5}$	$6,9064 \cdot 10^{-6}$	$5,3063 \cdot 10^{-6}$
6	$1,8925 \cdot 10^{-3}$	$5,3393 \cdot 10^{-5}$	$1,1764 \cdot 10^{-6}$	$5,0464 \cdot 10^{-7}$
7	$1,5110 \cdot 10^{-3}$	$4,6339 \cdot 10^{-5}$		$4,3806 \cdot 10^{-8}$
8	$1,8389 \cdot 10^{-3}$			
9	$1,8421 \cdot 10^{-3}$			
10	$7,6494 \cdot 10^{-4}$			

2. Numerical Results

In order to demonstrate how the numerical method works, we show two examples.

Example 1. As the first example, we solve the following equation:

$$\frac{1}{\Gamma(\beta)} \left(\int_0^{t/2} \frac{e^{(t+s)}x(s)}{(t-s)^\beta} ds + \int_{t/2}^{2t/3} \frac{tsx(s)}{(t-s)^\beta} ds + \int_{2t/3}^t \frac{(s+2)^3x(s)}{(t-s)^\beta} ds \right) = f(t),$$

where the right hand side of the equation is chosen such that the exact solution is $\bar{x}(t) = \sin(t)$. Calculations are performed for $t \in [0, T]$, the order of the fractional integration is $\beta = \frac{1}{2}$. Table 1 shows the errors $\varepsilon = \max_{0 \leq i \leq N} |\bar{x}(t_i) - X_m(t_i)|$ for different values of m and h .

Example 2. As the second example, we solve the equation

$$\frac{1}{\Gamma(\beta)} \left(\int_0^{\sin \frac{t}{2}} \frac{\sqrt{t+sx(s)}}{(t-s)^\beta} ds + \int_{\sin \frac{t}{2}}^{\sin \frac{2t}{3}} \frac{(t+s)x(s)}{(t-s)^\beta} ds + \int_{\sin \frac{2t}{3}}^t \frac{(t^2+s+3)x(s)}{(t-s)^\beta} ds \right) = f(t),$$

where, as well as in the previous example, the right hand side of the equation is chosen such that the exact solution is $\bar{x}(t) = \frac{\ln(t+1)t^2}{\sqrt{t+1}}$, $t \in [0, T]$ and $\beta = \frac{1}{2}$.

Table 2 shows the errors $\varepsilon = \max_{0 \leq i \leq N} |\bar{x}(t_i) - X_m(t_i)|$ for different values of m and h .

The examples show the following interesting feature: as the step decreases from $h = 10^{-4}$ to $h = 10^{-5}$, the error is significantly decreased at Example 1, and remains at the same level at Example 2. Note that the error in the approximate solution to the second model problem is somewhat worse than for the first one. This is due to less smooth components $K_i(t, s)$ and a more complex structure of the exact solution. In the future, we intend to develop more accurate quadratures for approximation integrals arising when calculating the coefficients of the system.

Table 2

Numerical results for Example 2

m	Errors			
	$h = 10^{-2}$	$h = 10^{-3}$	$h = 10^{-4}$	$h = 10^{-5}$
3	$6,2792 \cdot 10^{-3}$	$6,3200 \cdot 10^{-3}$	$6,3211 \cdot 10^{-3}$	$6,3212 \cdot 10^{-3}$
4	$1,6741 \cdot 10^{-3}$	$1,7364 \cdot 10^{-3}$	$1,7370 \cdot 10^{-3}$	$1,7371 \cdot 10^{-3}$
5	$4,9248 \cdot 10^{-4}$	$5,3815 \cdot 10^{-4}$	$5,3845 \cdot 10^{-4}$	$5,3846 \cdot 10^{-4}$
6	$3,1440 \cdot 10^{-4}$	$1,9342 \cdot 10^{-4}$	$1,9291 \cdot 10^{-4}$	$1,9292 \cdot 10^{-4}$
7	$3,1069 \cdot 10^{-4}$	$7,9240 \cdot 10^{-5}$	$8,0305 \cdot 10^{-5}$	$8,0298 \cdot 10^{-5}$
8		$3,8965 \cdot 10^{-5}$	$3,8303 \cdot 10^{-5}$	$3,8296 \cdot 10^{-5}$
9		$2,0569 \cdot 10^{-5}$	$2,0370 \cdot 10^{-5}$	$2,0406 \cdot 10^{-5}$
10		$1,0447 \cdot 10^{-5}$	$1,1771 \cdot 10^{-5}$	$1,1840 \cdot 10^{-5}$
11			$7,4383 \cdot 10^{-6}$	$7,3277 \cdot 10^{-6}$
12			$4,5928 \cdot 10^{-6}$	$4,7671 \cdot 10^{-6}$
13			$3,2945 \cdot 10^{-6}$	$3,2260 \cdot 10^{-6}$
14			$2,7023 \cdot 10^{-6}$	$2,2571 \cdot 10^{-6}$
15			$1,2863 \cdot 10^{-6}$	$1,6224 \cdot 10^{-6}$

Conclusion

In this article, we present the numerical collocation method that can be used to solve linear Volterra fractional order integral equations of the first kind with piecewise continuous kernel. The proposed quadrature formula effectively employ the restriction of the considered weakly regular equations with fractional order integrals and piecewise continuous kernels.

Acknowledgment. *Authors owe a special debt of gratitude to Professor Nikolai A. Sidorov for his generous, thoughtful support and for sharing his vast knowledge.*

This work was partially supported by the base part of the Government Assignment for Scientific Research from the Ministry of Science and Higher Education of Russia (project code: FZZS-2020-0039).

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Received October 5, 2020

ЧИСЛЕННЫЙ АНАЛИЗ ДРОБНЫХ ИНТЕГРАЛЬНЫХ ДИНАМИЧЕСКИХ МОДЕЛЕЙ С КУСОЧНО-НЕПРЕРЫВНЫМИ ЯДРАМИ*А.Н. Тында¹, Д.Н. Сидоров^{2,3}, И.Р. Муфтахов^{2,4}*¹Пензенский государственный университет, г. Пенза, Российская Федерация²Институт систем энергетики имени Л.А. Мелентьева СО РАН, г. Иркутск, Российская Федерация³Иркутский национальный исследовательский технический университет, г. Иркутск, Российская Федерация⁴Главный вычислительный центр ОАО «РЖД», г. Иркутск, Российская Федерация

Интегральные уравнения Вольтерра находят свое применение во многих областях, включая математическую физику, теорию управления, механику, электротехнику, и в различных отраслях промышленности. В частности, динамические модели Вольтерра с разрывными ядрами эффективно используются в энергетике для определения режимов работы накопителей энергии, а также для решения задачи выравнивания нагрузки. В статье предлагается численный метод решения линейных интегральных уравнений Вольтерра первого рода дробного порядка интегрирования с кусочно-непрерывными ядрами. Разработанный подход основан на методе полиномиальной коллокации и эффективно аппроксимирует такие слабо сингулярные интегралы. Эффективность предложенного численного метода иллюстрируется на двух примерах.

Ключевые слова: интегральные уравнения Вольтерра; численный метод; сходимость; сингулярность; дробный интеграл.

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Поступила в редакцию 5 октября 2020 г.